

$(\text{CF}_3)_2\text{PPF}_2 \cdot \text{BH}_3$  was effected when distillation through the  $-95^\circ$  trap was slow.

The vapor density of  $(\text{CF}_3)_2\text{PPF}_2 \cdot \text{BH}_3$  also corresponds well with theory (mol wt: calcd, 251.8; found, 253). However, the molecule dissociates easily as evidenced by the observation of appreciable amounts of  $\text{B}_2\text{H}_6$  and  $(\text{CF}_3)_2\text{PPF}_2$  in the infrared spectra and the lack of a molecular ion in the mass spectrum of  $(\text{CF}_3)_2\text{PPF}_2$  (Figure 1). Other easily dissociated species such as  $\text{B}_4\text{H}_8 \cdot \text{CO}$  also show only fragment ions by conventional mass spectrometry.<sup>42</sup>

**Instrumentation.**—Infrared spectra were obtained with an 87-mm gas cell equipped with CsI windows on either a Beckman IR-10 or a Beckman IR-12. The Raman samples were sealed

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in capillary tubes and the spectra determined on a SPEX Ramalog 1401 using an argon ion laser. A Consolidated Electro-dynamics Model 21-103B operating at 70 eV was used to obtain the mass spectra. The  $^1\text{H}$ ,  $^{19}\text{F}$ ,  $^{31}\text{P}$ , and  $^{11}\text{B}$  nmr spectra were determined with a Varian HR-100 operating at 100.0, 94.1, 40.5, and 32.1 MHz, respectively. A Cary 14 was used to obtain the uv spectra of samples in a 103-mm quartz gas cell.

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CONTRIBUTION FROM THE DEPARTMENT OF CHEMISTRY,  
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## The Phosphorus-Phosphorus Bond. III. The Microwave Spectrum, Conformation, and Dipole Moment of Phosphinodifluorophosphine<sup>1</sup>

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The rotational spectrum of phosphinodifluorophosphine was assigned with the following rotational constants (MHz):  $A = 6916.53$ ,  $B = 3372.94$ ,  $C = 2648.59$  for  $\text{H}_2\text{PPF}_2$ ;  $A = 6625.56$ ,  $B = 3192.44$ ,  $C = 2522.37$  for  $\text{D}_2\text{PPF}_2$ . The dipole moment was measured as  $1.71 \pm 0.01$  D, with principal axes components  $\mu_a = 1.25$  D and  $\mu_c = 1.16$  D. From the rotational constants and dipole moment data, the conformation was determined as trans. With the assumption that  $d(\text{PH}) = 1.42 \pm 0.01$  Å, the following structural parameters were derived:  $d(\text{PP}) = 2.218 \mp 0.038$  Å,  $d(\text{PF}) = 1.587 \pm 0.013$  Å,  $\angle\text{HPH} = 93.2 \mp 1.0^\circ$ ,  $\angle\text{HPP} = 90.3 \mp 0.4^\circ$ ,  $\angle\text{FPF} = 98.2 \mp 1.2^\circ$ , and  $\angle\text{FPP} = 97.2 \pm 1.6^\circ$ .

### Introduction

Structure investigations of diphosphines are of interest because of the possibility of conformational isomerism. Although no precise structural data are available, the infrared and Raman spectra have been interpreted in terms of a gauche configuration for  $\text{P}_2\text{H}_4$ <sup>2,3</sup> and a trans configuration for  $\text{P}_2\text{F}_4$ <sup>4</sup> and  $\text{P}_2\text{Cl}_4$ .<sup>5</sup> Vibrational spectra and X-ray data also confirm the trans symmetry for  $\text{P}_2\text{I}_4$ .<sup>6,7</sup> Vibrational spectra for  $(\text{CH}_3)_4\text{P}_2$  indicated a gauche-trans rotameric mixture in the liquid phase but only the trans form in the solid.<sup>8</sup> A recent electron-diffraction study of gaseous  $(\text{CH}_3)_4\text{P}_2$  indicated that the dihedral angle ( $\phi$ ) was  $164 \pm 23^\circ$ , measured from the cis configuration<sup>9</sup> ( $\phi = 0^\circ$  for cis,  $180^\circ$  for trans). Finally, semiempirical SCF and extended Hückel MO calculations for  $\text{P}_2\text{H}_4$  and  $\text{P}_2\text{F}_4$  suggested that the stability of the various conformations decreases in the order gauche > cis > trans.<sup>10</sup>

In the related compounds  $\text{F}_2\text{PN}(\text{CH}_3)_2$  and  $\text{F}_2\text{PNH}_2$ ,

a pyramidal arrangement of neighboring atoms about phosphorus was found; however, the arrangement about the nitrogen was planar and the  $\text{NX}_2$  plane bisected the FPF angle.<sup>11,12</sup>  $\text{PF}_2\text{PH}_2$  has been recently synthesized,<sup>1,13</sup> in view of the above data it was of interest to determine its conformation.

### Experimental Section

**Apparatus.**—The spectra were obtained with a Stark modulated spectrometer which has been described previously.<sup>14</sup> The transitions were measured with a precision of about 0.2 MHz. The absorption cell was cooled with Dry Ice ( $-78^\circ$ ) for all observations.

**Materials.**—The  $\text{H}_2\text{PPF}_2$  was prepared by the reaction of  $\text{PF}_2\text{I}$  and  $\text{PH}_3$  and identified by comparison with the reported infrared, nmr, and mass spectra.<sup>1,13</sup> The  $\text{D}_2\text{PPF}_2$  was made by substituting  $\text{PD}_3$  obtained by hydrolysis of calcium phosphide. The synthesis of  $\text{D}_2\text{PPF}_2$  was confirmed by infrared spectroscopy.

**Spectra.**—The transitions which were assigned for the two isotopic species are listed in Table I. The rotational constants listed in Table II were obtained by a least-squares fit of the observed transitions. The assignment was based on the Stark effect and the rigid-rotor fit for both species.

In addition to the transitions from  $\text{H}_2\text{PPF}_2$ , absorptions from the decomposition products  $\text{PF}_2\text{H}$  and  $\text{PF}_3$  were also noted. A weaker satellite line was found accompanying many of the ground-state lines. This was attributed to a vibrational satellite since the lowest fundamental has been assigned at 209

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TABLE I  
OBSERVED AND CALCULATED TRANSITIONS (MHz)  
FOR H<sub>2</sub>PPF<sub>2</sub> AND D<sub>2</sub>PPF<sub>2</sub>

Transition	H <sub>2</sub> PPF <sub>2</sub>		D <sub>2</sub> PPF <sub>2</sub>	
	$\nu(\text{obsd})$	$\nu(\text{obsd}) - \nu(\text{calcd})$	$\nu(\text{obsd})$	$\nu(\text{obsd}) - \nu(\text{calcd})$
1 <sub>10</sub> → 2 <sub>20</sub>	23,498.61	0.33	22,488.05	0.14
1 <sub>11</sub> → 2 <sub>21</sub>	24,122.66	0.14	23,069.03	-0.11
2 <sub>11</sub> → 3 <sub>21</sub>	29,183.96	0.12	27,878.03	-0.10
2 <sub>11</sub> → 3 <sub>12</sub>	19,083.84	0.33		
2 <sub>02</sub> → 3 <sub>12</sub>	24,176.33	0.40	22,952.08	0.13
2 <sub>02</sub> → 3 <sub>03</sub>	17,676.40	0.24		
3 <sub>03</sub> → 4 <sub>04</sub>	23,182.31	0.02	22,050.95	-0.33
3 <sub>13</sub> → 4 <sub>14</sub>	22,460.91	0.12	21,361.10	-0.50
3 <sub>12</sub> → 4 <sub>13</sub>	25,306.54	0.20	23,997.77	0.01
3 <sub>22</sub> → 4 <sub>23</sub>	24,007.90	0.03	22,789.50	-0.37
3 <sub>21</sub> → 4 <sub>22</sub>	24,910.92	-0.13	23,597.30	-0.05
4 <sub>04</sub> → 5 <sub>05</sub>	28,505.98	-0.20	27,132.49	0.21
4 <sub>14</sub> → 5 <sub>15</sub>	27,939.98	-0.02	26,579.98	0.28
4 <sub>13</sub> → 5 <sub>14</sub>	31,385.29	-0.01	29,779.53	0.04

TABLE II  
ROTATIONAL CONSTANTS (MHz), MOMENTS OF  
INERTIA (AMU Å<sup>2</sup>), AND SECOND MOMENTS  
(AMU Å<sup>2</sup>) FOR H<sub>2</sub>PPF<sub>2</sub> AND D<sub>2</sub>PPF<sub>2</sub>

	H <sub>2</sub> PPF <sub>2</sub>	D <sub>2</sub> PPF <sub>2</sub>
A	6916.53	6625.56
B	3372.94	3192.44
C	2648.59	2522.37
I <sub>a</sub>	73.0902	76.3001
I <sub>b</sub>	149.8784	158.3525
I <sub>c</sub>	190.8679	200.4190
P <sub>a</sub>	133.8280	141.2357
P <sub>b</sub>	57.0398	59.1833
P <sub>c</sub>	16.0503	17.1168

TABLE III  
OBSERVED AND CALCULATED STARK COEFFICIENTS (10<sup>6</sup> × MHz cm<sup>2</sup>/V<sup>2</sup>)  
AND THE CALCULATED SECOND-ORDER PERTURBATION COEFFICIENTS

Transition	Obsd	Calcd			K <sub>a</sub>	K <sub>b</sub>	K <sub>c</sub>
		I <sup>a</sup>	II <sup>a</sup>				
2 <sub>02</sub> → 3 <sub>12</sub>	m = 1	1.062	1.062	1.062	-0.0847	0.4127	0.7923
2 <sub>11</sub> → 3 <sub>12</sub>	m = 1	-0.3476	-0.3493	-0.3504	-0.2650	-0.1503	0.0452
2 <sub>11</sub> → 3 <sub>12</sub>	m = 2	-1.250	-1.250	-1.250	-0.9742	-0.8838	0.1950
2 <sub>02</sub> → 3 <sub>03</sub>	m = 0	-0.1548	-0.1531	-0.1544	-0.0990	-0.0004	-0.0004
3 <sub>22</sub> → 4 <sub>23</sub>	m = 0	0.0823	0.0823	0.0823	0.0218	0.0427	0.0362

<sup>a</sup> Consult text for definition.

cm<sup>-1</sup>,<sup>1</sup> it was not extensively investigated. Attempts were made to identify additional low *J* lines arising from a second conformation but none was found indicating that another isomeric form is markedly less abundant. The infrared, Raman, and nmr spectra also have offered no evidence suggesting a second rotamer.<sup>1,13</sup>

**Dipole Moment.**—The Stark shift measurements were made with a precision dc voltage source (Fluke 413B) using the method described by Kirchhoff.<sup>15</sup> The second-order Stark coefficients are listed in Table III. The data were consistent to better than 1%. The effective guide spacing was determined using  $\mu = 0.7152$  D for OCS.<sup>16</sup>

### Analysis

**Selection Rules and Dipole Components.**—The symmetry of the conformer can be determined in principle by noting the projections that the total dipole moment makes along the principal axes. For a species with no symmetry such as *gauche*, the total dipole will have nonzero components for  $\mu_a$ ,  $\mu_b$ , and  $\mu_c$ . For a species such as *trans* with *C<sub>s</sub>* symmetry, there will be two components since the total dipole moment must lie in the symmetry plane. For a species with a *C<sub>2</sub>* axis such as a

planar *C<sub>2v</sub>* conformation, only one component will be nonzero. The existence of the dipole components can be determined by the types of transitions observed and by the Stark effect.

The transitions that were assigned in Table I consisted of both  $\mu_a$  and  $\mu_c$  types of about equal intensity. The  $\mu_b$  type transitions were also predicted and a search was made for them. Although a few weak transitions were found at these predicted frequencies, most of the intense lines could not be found and hence no support for *b* type selection rules was obtained. Since the intensity of a transition is proportional to the dipole component squared,  $\mu_b$  was estimated to be  $\leq 0.1$  D.

The Stark effect was also quantitatively investigated in order to place a limit on the magnitude of  $\mu_b$ , since the frequency shifts ( $\Delta\nu$ ) in the electric field ( $\epsilon$ ) are determined by

$$\Delta\nu/\epsilon^2 = K_a\mu_a^2 + K_b\mu_b^2 + K_c\mu_c^2 \quad (1)$$

$K_a$ ,  $K_b$ , and  $K_c$  are the second-order perturbation coefficients. These are readily calculated if the rotational constants are known.

The experimental Stark coefficients ( $\Delta\nu/\epsilon^2$ ) which are listed in Table III were fit by a least-squares procedure for the three dipole components giving  $\mu_a^2 = 1.54$  D,  $\mu_b^2 = 0.013$  D, and  $\mu_c^2 = 1.35$  D. The fit for these values is listed under "Calcd I." If  $\mu_b^2$  was fixed identical with zero and the least-squares procedure was repeated, then the values obtained were  $\mu_a^2 = 1.55$  D

and  $\mu_c^2 = 1.36$  D. The fit for these values is listed under "Calcd II." Both fits are equally good and are within experimental uncertainty; hence it appears that  $\mu_b^2 \leq 0.01$  D. Since the isotope shifts discussed below show that  $\mu_b = 0$ , we have chosen the second set of dipole components to calculate the total dipole moment. They give  $|\mu_a| = 1.25$  D,  $|\mu_c| = 1.16$  D, and a total dipole moment of  $1.71 \pm 0.01$  D. The values of  $K_a$ , etc., have also been listed in Table III to show the sensitivity of the equations to the dipole components.

The transitions in Table III were selected to ensure that the experimental data could be safely interpreted using second-order perturbation theory. This was necessary since calculations showed that contributions from a possible small value of  $\mu_b$  in second-order would be comparable to fourth-order corrections to the Stark energy from  $\mu_a$  and  $\mu_c$ . Ignoring these higher order corrections could lead to an erroneous conclusion concerning  $\mu_b$ .<sup>17</sup> The transitions in Table III were selected after comparing the Stark shifts calculated from preliminary values of  $\mu_a$  and  $\mu_c$  using eq 1 with

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those calculated by a diagonalization of the Stark energy matrix<sup>18</sup> for  $J$  up to 6. Only transitions giving the same Stark shift were retained in Table III. This procedure neglected higher order corrections from  $\mu_b$ , which upon examination of the form of the fourth-order corrections derived by Kirchhoff<sup>15</sup> is justified if  $\mu_b$  is small.

In summary, the absence of evidence for  $\mu_b$  type selection rules indicates that it is certainly small and is probably zero. This strongly suggests that the conformer is not gauche, although it cannot be rigorously eliminated since  $\mu_b$  could accidentally be close to zero. From an examination of bond dipole arguments, this possibility seems unlikely. Since  $\mu_a$  and  $\mu_c$  are both nonzero, structures possessing a  $C_2$  axis are also eliminated. The dipole data are most consistent with a conformer having  $C_s$  symmetry.

**Isotope Shift.**—The moments of inertia for the normal isotopic species do not uniquely determine the symmetry because of the large number of structural parameters. However, the additional data from the  $D_2PPF_2$  species do help to define the symmetry by the following arguments.

(1) The second moments of inertia (Table II) for both isotopic species [ $P_a = 1/2(I_b + I_c - I_a)$ , etc.] show that  $P_a$ ,  $P_b$ , and  $P_c$  changed upon deuteration. This indicates that the conformer cannot have a plane of symmetry containing two hydrogens, thus eliminating a structure analogous to  $H_2NPF_2$ .

(2) The differences in moments of inertia [ $I_a$  (normal) -  $I_a$  (deuterium)] expected for the trans, cis, and gauche conformations based on plausible trial structures are listed in Table IV. Comparing the isotopic

TABLE IV  
CHANGES IN MOMENTS OF INERTIA (AMU Å<sup>2</sup>) UPON  
DEUTERATION PREDICTED FOR SEVERAL CONFORMATIONS  
BASED ON AN INITIALLY ASSUMED STRUCTURE

	$\Delta I_a$	$\Delta I_b$	$\Delta I_c$
Trans	3.18	9.11	10.37
Gauche	5.72	6.58	8.94
Cis	4.12	6.32	6.63
Obsd	3.21	8.47	9.55

shifts with those experimentally measured shows that the trans form is favored. The trial structure employed bond lengths and angles from  $PH_3$  and  $PF_3$  and a PP bond of 2.22 Å; it was used initially to identify the spectra of both isotopic species.

(3) Kraitichman's equations for double substitution<sup>19</sup> in a molecule with  $C_s$  symmetry gave hydrogen coordinates of  $|a| = 1.928$ ,  $|b| = 1.032$ , and  $|c| = 0.758$  Å. These were within 1-5% of those predicted by our trial trans structure. If the isomer was cis, the hydrogen  $a$  and  $c$  coordinates should have been considerably smaller since the hydrogens would be closer to the center of mass. If the hydrogen coordinates were calculated assuming single rather than double deuterium substitution, the coordinates obtained were unreasonably large precluding the possibility that the wrong isotopic species was assigned.

In conclusion the changes in the moments of inertia upon deuteration eliminate the gauche or cis conforma-

tion or a structure like that found for  $H_2NPF_2$ . They are most consistent with the trans conformation.

**Structural Parameters.**—There are seven structural parameters for a trans conformation; the data from the two species cannot fix them without one assumption. These structural parameters were calculated for assumed PH bond lengths between 1.38 and 1.44 Å, a range encompassing all known PH bond lengths. The results are included in Table V. It can be seen that

TABLE V  
SETS OF STRUCTURAL PARAMETERS WHICH FIT THE  
OBSERVED ROTATIONAL CONSTANTS (BOND  
LENGTHS IN ÅNGSTRÖMS, BOND ANGLES IN  
DEGREES)

$d(PH)$	$d(PF)$	$d(PP)$	$\angle HPH$	$\angle HPP$	$\angle FPF$	$\angle FPP$
1.380	1.535	2.365	96.8	92.1	103.0	89.9
1.390	1.548	2.325	95.8	91.6	101.4	91.8
1.400	1.558	2.290	94.9	91.2	100.8	93.8
1.410	1.574	2.251	94.1	90.8	99.4	95.6
1.420 <sup>a</sup>	1.587	2.218	93.2	90.3	98.2	97.2
1.430	1.604	2.184	92.4	89.9	97.0	98.6
1.440	1.623	2.142	91.6	89.5	95.4	100.2

<sup>a</sup> The structures associated with  $d(PH) = 1.42 \pm 0.01$  Å are preferred by the authors. See text.

some parameters such as  $d(PP)$  and  $d(PF)$  vary greatly with  $d(PH)$ . One parameter  $\angle PPH$  is particularly insensitive to the value assumed for  $d(PH)$ .

It seems justified to limit the range of possible structures to those associated with the assumption  $d(PH) = 1.42 \pm 0.01$  Å. The reasons are (a)  $d(PF)$  is expected to be larger than 1.57 Å based on  $PF_3$  (1.570 Å)<sup>20</sup> and  $PHF_2$  (1.582 Å),<sup>14</sup> (b)  $\angle FPF$  is expected to be between 97 and 100° based on  $PF_3$  (97.8°) and  $HPF_2$  (99°), (c) most PP single bond lengths fall in the range  $2.22 \pm 0.04$  Å,<sup>9</sup> and (d)  $d(PH)$  is expected to be near 1.42 Å based on  $PH_3$  (1.419 Å)<sup>21</sup> and  $PHF_2$  (1.412 Å). Hence, the structural parameters most consistent with the experimental data and the results from other molecules are  $d(PH) = 1.42 \pm 0.01$  Å,  $d(PP) = 2.18 \mp 0.038$  Å,  $d(PF) = 1.587 \pm 0.013$  Å,  $\angle HPH = 93.2 \mp 1^\circ$ ,  $\angle HPP = 90.3 \mp 0.4^\circ$ ,  $\angle FPF = 98.2 \mp 1.2^\circ$ , and  $\angle FPP = 97.2 \pm 1.6^\circ$ . The uncertainties encompass the range of parameters based on the assumed PH bond length of  $1.42 \pm 0.01$  Å.

## Discussion

This section will center upon the two conclusive results of this study, *viz.*, the conformation and the dipole moment. The structural parameters do not warrant a detailed analysis due to their large uncertainties.

It requires several assumptions in order to analyze the total dipole moment of  $H_2PPF_2$  with a bond dipole moment model particularly because of the contributions of the lone pairs. Nevertheless it is interesting to speculate on the size of the PP bond moment by transferring PF and PH bond moments from other molecules and assuming that any lone-pair contributions are contained in those bond moments. Using the bond dipole moments from  $PF_3$ <sup>22</sup> and  $PH_3$ <sup>23</sup> and the structure discussed above gave  $\mu_a = 0.79$  D (observed 1.25 D) and  $\mu_c = 1.32$  D (observed 1.16 D). A PP

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bond moment of 0.5 D will nicely raise  $\mu_a$  and lower  $\mu_c$  to the observed values only if the negative end is directed to the phosphorus with the fluorine substituents. This value suggests an inductive effect from the fluorines, and the direction of the bond dipole is consistent with the greater Lewis basicity toward borane of the phosphorus with the attached fluorines. It is interesting to note that the same value for the PP bond moment fits both  $\mu_a$  and  $\mu_c$  for  $H_2PPF_2$ , implying that the bond dipole moments in  $PF_3$ ,  $PH_3$ , and  $H_2PPF_2$  are internally consistent.

This investigation shows that the stable conformation is trans with no evidence for appreciable amounts of a gauche conformer. This is particularly interesting for two reasons. First, there appears to be no tendency for the lone pair on the  $PH_2$  group to form a  $d_{\pi}-p_{\pi}$  bond to the  $PF_2$  end since this would make the  $PH_2$  group planar. This effect can be invoked to rationalize the planar arrangement around the nitrogen in  $H_2NPF_2$ <sup>12</sup> and  $(CH_3)_2NPF_2$ .<sup>11</sup> An analogy to this result has been observed with  $N(SiH_3)_3$  which is planar<sup>24</sup> while  $P(SiH_3)_3$  is pyramidal.<sup>25</sup>

Second, the stability of the trans form is interesting in view of the electrical asymmetry in  $H_2PPF_2$ ; apparently an internal attraction between the hydrogens and

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fluorines is not influential enough to stabilize the gauche isomer. There is, however, some suggestion of an interaction involving the hydrogens since the PPH angle is a "small"  $90^\circ$ . Since  $H_2PPF_2$  readily decomposes to  $HPF_2$ , it seems preferable to call the interaction an attraction presumably toward the lone pair on the adjacent phosphorus.

The data now available indicate that all gaseous diphosphines have the trans conformation except  $H_4P_2$ . A trans conformation seems readily justified invoking steric repulsion and/or lone pair interactions. These two effects should be the smallest in  $H_4P_2$  so a gauche conformation would be more likely; nevertheless some stabilizing interaction must still operate to give the gauche conformation. The results of this investigation certainly suggest that further studies of the molecular and electronic structure of  $H_4P_2$  and its potential function for rotation about the PP bond are warranted.

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## Vibrational Spectra and Assignments for Isotopic Species of Tetraborane(10)<sup>1</sup>

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Four isotopically enriched species of tetraborane(10) have been prepared involving hydrogen, deuterium, boron-10, and boron-11. The infrared spectra of these compounds have been investigated in the gas and solid states in the spectral region from 400 to 4000  $cm^{-1}$ . Low-resolution Raman spectra in the liquid and solid states have been obtained and qualitative polarization measurements made. Vibrational modes of this molecule have been described on the basis of a symmetry analysis and a group vibration approximation. Tentative vibrational assignments have been made for 34 of the 36 fundamentals.

### Introduction

Despite the extensive chemical interest in the boron hydrides and the intriguing nature of the hydrogen bridge bond which these compounds contain, spectroscopic attention has been focused almost entirely upon the simplest member of the series, diborane. Experimental difficulties caused by the reactive nature and general thermal instability of the higher hydrides, together with spectroscopic complications inherent in their relatively large size and the lack of isotopic purity of natural boron, seem to have been effective deterrents to study. These compounds clearly deserve considerably more attention than they have been accorded.

In the present work, the vibrational spectra of four

isotopic varieties of tetraborane (10),  $B_4H_{10}$ , have been investigated in the gas, liquid, and solid states under low resolution in order to arrive at a preliminary set of vibrational assignments. The infrared spectrum of this compound prepared from boron containing the natural isotopic abundance has been reported in the literature, the most detailed results being those of Spielman and Burg,<sup>3</sup> who published the infrared spectrum of the vapor for both the hydrogen and deuterium compounds. Qualitative assignments were given based on frequency shifts and analogies with the spectra of other compounds. No Raman data have been found. A few infrared frequencies for the mixed isotopic species,  $B_4H_8D_2$ , have also been reported.<sup>4</sup> The only other higher boron hydride which has been given any signifi-

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